Introduction to Ordinary Differential Equations

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Outline

- What is a differential equation?
- Initial Value Problems
 - Linear first order differential equations
 - Second order differential equations
 - Recasting high order differential equations as a system of first order differential equations
- Boundary Value Problems
- Solution techniques for nonlinear differential equations
 - Power series solutions
 - Perturbation theory concept
- Oncluding Remarks

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Population growth is commonly modeled with differential equations. In the following equation: t = time, P = population and k = proportionality constant. k represents the constant ratio between the growth rate of the population and the size of the population.

$$\frac{dP}{dt} = kP$$

In this particular equation, the left hand side represents the growth rate of the population being proportional to the size of the population P. This is a very simple example of a first order, ordinary differential equation. The equation only contains first order derivatives and there are no partial derivatives.

An initial value problem consists of a differential equation and an initial condition. So, going back to the population example, the following is an example of an initial value problem:

$$\frac{dP}{dt} = kP, P(0) = P_0$$

The solution to this set of equations is a function, call it P(t), that satisfies both equations.

• The standard form for a first-order differential equation is

$$\frac{dy}{dt} = f(t, y)$$

where the right hand side represents the function f that depends on the independent variable, t, and the dependent variable, y.

General Solutions to a Differential Equation

Let's look at a simple example and walk through the steps of finding a general solution to the following equation

$$\frac{dy}{dt} = (ty)^2$$

We will simply "separate" write as "separate" the variables then integrate the both sides of the equation to find the general solution.

$$\frac{dy}{dt} = t^2 y^2$$
$$\frac{1}{y^2} dy = t^2 dt$$
$$\int \frac{1}{y^2} dy = \int t^2 dt$$

$$-y^{-1} = \frac{t^3}{3} + c$$
$$-\frac{1}{y} = \frac{t^3}{3} + c$$
$$y = -\frac{1}{\frac{t^3}{3} + c}$$
$$\Rightarrow y(t) = -\frac{3}{t^3 + c_1}$$

where c_1 is any real number.

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Linear First Order Differential Equations

Initial value problems consist of a differential equation and an initial value. We will work through the example below:

$$rac{dx}{dt} = -xt;$$
 $x(0) = rac{1}{\sqrt{\pi}}$

First we will need to find the general solution to $\frac{dx}{dt} = -xt$, then use the initial value $x(0) = \frac{1}{\sqrt{\pi}}$ to solve for *c*. Since we do not know what x(t) is, we will need to "separate" the equation before integrating.

$$\frac{dx}{dt} = -xt$$
$$-\frac{1}{x}dx = t dt$$
$$\int -\frac{1}{x}dx = \int t dt$$

Linear First Order Differential Equations Continued

$$-\ln x = \frac{t^2}{2} + c$$

$$x = e^{-(\frac{t^2}{2} + c)}$$

$$x = e^{-(\frac{t^2}{2})}e^{-c}$$

$$x = ke^{-\frac{t^2}{2}}$$

The above function of t is the general solution to $\frac{dx}{dt} = -xt$ where k is some constant. Since we have the initial value $x(0) = \frac{1}{\sqrt{\pi}}$, we can solve for k.

Thus we can see that the solution to the initial value problem

$$\frac{dx}{dt} = -xt; x(0) = \frac{1}{\sqrt{\pi}}$$

is

$$x(0) = \frac{1}{\sqrt{\pi}} = ke^{-\frac{0^2}{2}}$$

$$x(t) = \frac{1}{\sqrt{\pi}} e^{-\frac{t^2}{2}}$$

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Image: Image:

Let's verify that this solution is correct. We will need to show

$$\begin{aligned} \frac{dx}{dt} &= x'(t) = f(t, x(t)) \\ \frac{dx}{dt} &= \frac{d}{dt} \left(\frac{1}{\sqrt{\pi}} e^{-\frac{t^2}{2}} \right) = \frac{1}{\sqrt{\pi}} e^{-\frac{t^2}{2}} \\ &\Rightarrow \frac{d}{dt} \left(\frac{1}{\sqrt{\pi}} e^{-\frac{t^2}{2}} \right) = -\frac{1}{\sqrt{\pi}} e^{-\frac{t^2}{2}} \end{aligned}$$

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Second order differential equations simply have a second derivative of the dependent variable. The following is a common example that models a simple harmonic oscillator:

$$\frac{d^2y}{dt^2} + \frac{k}{m}y = 0$$

where m and k are determined by the mass and spring involved. This second order differential equation can be rewritten as the following first order differential equation:

$$\frac{dv}{dt} = -\frac{k}{m}y$$

where v denotes velocity.

Second Order Differential Equations Continued

Referring back to some calculus knowledge, if v(t) is velocity, then $v = \frac{dy}{dt}$. Thus, we can substitute in $\frac{dv}{dt}$ into our second order differential equation and essentially turn it into a first order differential equation.

$$\frac{d^2y}{dt^2} = -\frac{k}{m}y \Leftrightarrow \frac{dv}{dt} = -\frac{k}{m}y$$

Now we have the following system of first order differential equations to describe the original second order differential equation:

$$\frac{dy}{dt} = v$$
$$\frac{dv}{dt} = -\frac{k}{m}y$$

Consider the following initial value problem:

$$\frac{d^2y}{dt^2} + y = 0$$

with y(0) = 0 and y'(0) = v(0) = 1. Let's show that y(t) = sin(t) is a solution. Let $v = \frac{dy}{dt}$, then we have the following system:

$$\frac{dy}{dt} = v$$
$$\frac{dv}{dt} = -y$$

Second Order Differential Equations Continued

$$\frac{dy}{dt} = \frac{d}{dt}\sin(t) = \cos(t) = v$$
$$\frac{dv}{dt} = -\sin(t) = -y$$
$$\Rightarrow \frac{d^2y}{dt^2} = -\sin(t)$$
$$\Rightarrow \frac{d^2y}{dt^2} + y = \frac{d^2(\sin(t))}{dt^2} + \sin(t)$$
$$= -\sin(t) + \sin(t) = 0$$

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High Order Differential Equations as a System

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Boundary Value Problems: The Basics

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Power Series Solutions

To demonstrate how to use power series to solve a nonlinear differential equation we will look at Hermite's Equation:

$$\frac{d^2y}{dt^2} - 2t\frac{dy}{dt} + 2py = 0$$

We will use the following power series and its first and second derivatives to make a guess:

$$y(t) = a_0 + a_1 t + a_2 t^2 + a_3 t^3 + \dots = \sum_{n=0}^{\infty} a_n t^n$$
(1)

$$\frac{dy}{dt} = a_1 + 2a_2t + 3a_3t^2 + 4a_4t^3 + \dots = \sum_{n=1}^{n} na_nt^{n-1}$$
(2)

$$\frac{d^2y}{dt^2} = 2a_2 + 6a_3t + 12a_4t^2 + \dots = \sum_{n=2}^{\infty} n(n-1)a_nt^{n-2}$$
(3)

From the previous equations we can conclude that

$$\begin{array}{rcl} y(0) & = & a_0 \\ y'(0) & = & a_1 \end{array}$$

Next we will substitute (1), (2) and (3) into Hermite's Equation and collect like terms.

$$\begin{aligned} \frac{d^2y}{dt^2} - 2t\frac{dy}{dt} + 2py &= 0 = (2a_2 + 6a_3t + 12a_4t^2 + \dots) \\ &- 2t(a_1 + 2a_2t + 3a_3t^2 + 4a_4t^3 + \dots) \\ &+ 2p(a_0 + a_1t + a_2t^2 + a_3t^3 + \dots) \\ &\Rightarrow (2pa_0 + 2a_2) + (2pa_1 - 2a_1 + 6a_3)t + \\ &(2pa_2 - 4a_2 + 12a_4)t^2 + (2pa_3 - 6a_3 + 20a_5)t^3 = 0 \end{aligned}$$

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Then from here, we will set all coefficients equal to 0 since the equation is equal to 0 and $t \neq 0$. We get the following sequence of equations:

$$2pa_0 + 2a_2 = 0$$

$$2pa_1 - 2a_1 + 6a_3 = 0$$

$$2pa_2 - 4a_2 + 12a_4 = 0$$

$$2pa_3 - 6a_3 + 20a_5 = 0$$

Then will several substitutions we arrive at the following set of equations:

$$\Rightarrow a_{2} = -pa_{0}$$

$$a_{3} = -\frac{p-1}{3}a_{1}$$

$$a_{4} = -\frac{p-2}{6}a_{2} = \frac{(p-2)p}{6}a_{0}$$

$$a_{5} = -\frac{p-3}{10}a_{3} = \frac{(p-3)(p-1)}{30}a_{1}$$

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Perturbation theory is used when a mathematical equation involves a small perturbation, usually ϵ . From here we create y(x) such that it is an expansion in terms of ϵ . For example

$$y(x) = y_0(x) + \epsilon y_1(x) + \epsilon^2 y_2(x) + \cdots$$

This summation is called a perturbation series and it has a nice feature that allows each y_i to be solved using the previous y_i 's. Consider the equation,

$$x^2 + x + 6\epsilon = 0, \quad \epsilon \ll 1 \tag{4}$$

Let's consider using perturbation theory to determine approximations for the roots of Equation (4).

Notice this equation is a perturbation of $x^2 + x = 0$. Let $x(\epsilon) = \sum_{n=0}^{\infty} a_n \epsilon^n$. This series will be substituted into (4) and powers of ϵ will be collected. Next we will calculate the first term of the series by setting $\epsilon = 0$ in (4). So the leading order equation is

$$a_0^2 + a_0 = 0 (5)$$

with solutions x = -1, 0. Thus $x(0) = a_0 = -1, 0$. Now the perturbation series are as follows

$$= 1 - a_1 \epsilon - a_2 \epsilon^2 - a_1 \epsilon + a_1^2 \epsilon^2 + a_1 a_2 \epsilon^3 - a_2 \epsilon^2 + a_1 a_2 \epsilon^3 + a_2^2 \epsilon^4 - 1 + a_1 \epsilon \\ = (1 - 1) + (-2a_1 + a_1 + 6)\epsilon + (-2a_2 + a_1^2 + a_2)\epsilon^2 + \mathcal{O}(\epsilon^3)$$

Perturbation Theory Concept Continued

$$x_1(\epsilon) = -1 + a_1\epsilon + a_2\epsilon^2 + \mathcal{O}(\epsilon^3)$$
(9)

and

$$x_2(\epsilon) = 0 + b_1\epsilon + b_2\epsilon^2 + \mathcal{O}(\epsilon^3)$$
(10)

Next, we will substitute in (9) into (4) while ignoring powers of ϵ greater than 2. Since we are only approximating the solution to the second-order, we can disregard the powers of ϵ greater than 2.

$$x^{2} + x + 6\epsilon = (-1 + a_{1}\epsilon + a_{2}\epsilon^{2})^{2} + (-1 + a_{1}\epsilon + a_{2}\epsilon^{2}) + 6\epsilon$$
$$\Rightarrow (-a_{1} + 6)\epsilon + (-a_{2} + a_{1}^{2})\epsilon^{2} + \mathcal{O}(\epsilon^{3})$$

(8)

Perturbation Theory Concept Continued

From here we take the coefficient of each power of ϵ and set it equal to zero. This step is justified because (4) is equal to zero and $\epsilon \neq 0$ so each coefficient must be equal to zero. Thus we have the following equations

$$\mathcal{O}(\epsilon^1): -a_1 + 6 = 0$$

 $\mathcal{O}(\epsilon^2): a_1^2 - a_2 = 0$

These equations will be solved sequentially. The results are $a_1 = 6$ and $a_2 = 36$. Thus the perturbation expansion for the root $x_1 = -1$ is:

$$x_1(\epsilon) = -1 + 6\epsilon + 36\epsilon^2 + \mathcal{O}(\epsilon^3)$$

The same process can be repeated for x_2 with the perturbation expansion for the root $x_2 = 0$ resulting in

$$x_2(\epsilon) = -6\epsilon - 36\epsilon^2 + \mathcal{O}(\epsilon^3)$$

Concluding Remarks

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Questions?

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